Geometric modeling of 3D woven preforms in composite T-joints

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Abstract
A common method to fabricate net-shaped three-dimensional (3D) woven preforms for composite T-joints is to weave flat 3D preforms via a standard weaving machine with variation in binder yarn path and then separate the preform in the form of a bifurcation. Folding introduces fiber architecture deformation at the 3D woven bifurcation area. In this paper, a geometric modeling approach is proposed to represent the realistic fiber architecture, as a preprocessor for finite element analyses to predict composite structural performance. Supported by X-ray micro-computed tomography (μCT), three important deformation mechanisms are observed including yarn stack shifting, cross-section bending, and cross-section flattening resulting from the folding process. Furthermore, a set of mathematical formulae for simulation of the deformations in the junction region are developed and satisfactory agreement is observed when compared with μCT scan results.

Keywords
T-joint, 3D woven fabrics, geometric modeling, deformation

Three-dimensional (3D) woven composites are reinforced by 3D woven preforms containing yarns in $x$, $y$, and $z$ directions. 3D woven preforms, specifically for composite applications, are generally classified into orthogonal, through-the-thickness angle interlock, and layer-to-layer angle interlock patterns. Design and manufacturing of emerging multi-axis 3D woven preforms incorporating bias yarn layers has also been proposed but the development is still at an early stage for applications.¹ 3D woven composites have drawn great attention in recent decades for their advantages such as higher through-thickness properties that can help to overcome the problem of delamination encountered by two-dimensional (2D) laminates.¹ The exploration of 3D reinforced composites has also been extended to load-bearing profiles like T-, I-, and Pi-shaped joints. There are three basic ways to manufacture 3D reinforced composite joints with woven preforms. The first method is to use 2D woven fabric layups, reinforcing the through-thickness direction with stitching or Z-pinning.²,³ The second way is to weave flat 3D preforms via a standard weaving machine with variation of binder yarn path to form locally unreinforced planes (bifurcations) and then fold the preforms into the desired shape.⁴⁻⁷ The third way is to fabricate the preform directly into a complex 3D shape through a specially developed 3D weaving loom.⁸⁻¹⁰ Although composites reinforced by the first method can offer better damage tolerance in terms of delamination than 2D laminated joints subjected to equivalent loads,²,¹¹ they can be less competitive in in-plane properties because of geometric defects or fiber damage caused by retrospectively inserting through-thickness reinforcements.¹² In addition, due to the high investment and innovation needed for developing a special 3D loom for the third kind of preform, 3D preforms woven by a conventional loom are popular for current composite T-joints.

However, deformation is likely to occur when folding a flat woven preform into a 3D shape. Because the internal architecture of a woven preform is complex, its
variability in geometry can directly affect material properties such as permeability and formability. At the micro-scale, the fiber misalignment in fiber bundles was also found to have an influence on many aspects of composite behaviour. Previous studies have also shown that reinforcement geometry plays an important role in determining the composite’s mechanical properties, as two types of 3D woven composite T-joints with only weave variation at the 3D woven bifurcation area (in which the formed resin pocket is usually referred to as the “noodle”) were tested under the same tensile loading case but very different load–displacement responses were observed (Figure 1(a)). Thus, in order to closely predict the properties of composite T-joints reinforced by a flat woven 3D preform, it is necessary to model the fiber architecture as realistically as possible.

Characterizing the realistic fiber architecture of 3D woven composites is difficult because of the complexity of the internal geometry, but this can benefit from the use of X-ray micro-computed tomography (μCT). Some general frameworks for geometric modeling of woven preforms with respect to characterization of yarn path, crimp, and cross-section shape have been provided in the literature. Furthermore, compaction-induced yarn crimp and cross-section shape change were studied with a number of models; however, as they are based on beam theory, numerically solved by the finite element method, they are computationally intensive. Some mathematical methods were also proposed to model yarn cross-section shape under compression. An energy-based method used to characterize yarn deformation in compression was developed by Lomov and Verpoest, but the method asks for a measurement of the mechanical properties of yarns in bending and compression. Chen et al. modeled yarn cross-section deformation of both single and multi-layer woven preforms under compaction at the mesoscopic level based on the assumption that yarn width deformation is negligible compared with the deformation of yarn thickness. The model is not applicable when yarn width change...
is significant, for instance for a yarn subjected to bending load. Cornelissen and Akkerman observed yarn cross-section flattening in a cantilever experiment of yarns bending under their own weight.\textsuperscript{28} They concluded that cross-section flattening would affect the yarn flexural rigidity by one order of magnitude but they were not able to derive a trend from the experiment. However, yarn cross-section deformation in bending is significant as observed in this study via $\mu$CT analysis and an effective model is needed to capture this feature.

This paper focuses on the geometric modeling of 3D woven preform deformation caused by folding a flat woven preform into a T-piece, which is a typical fabrication process for composite T-joints. Through the $\mu$CT analysis of a compacted preform, three important deformation mechanisms specifically for this type of preform were observed, namely yarn stack shifting, cross-section bending and cross-section flattening resulting from the folding process. These features have been geometrically modeled and the models have been validated against $\mu$CT analysis data of a 3D woven T-joint preform.

### Material characterization

The example preform reinforcing the composite T-joints is a 3D orthogonal weave provided by Sigmatex, UK, based on Hexcel IM7 12K carbon fiber, which is woven flat on a Jacquard machine and folded into a T shape. 3D orthogonal weaves are 3D structures containing straight warp and weft yarns in the $x$, $y$ directions and binders (also called “warp weavers”) that interlace with other yarns, at times, orientated in the through-thickness direction, providing high stiffness under tensile loading due to low yarn waviness.\textsuperscript{10} The preform consists of eight layers of warp yarns and nine layers of weft yarns in the web, and four layers of warp yarns and five layers of weft yarns in the flange. A schematic woven structure for this preform before and after bifurcation is shown as “3D woven type 1” in Figure 1(b).

To characterize the intrinsic geometric features of this type of preform for composite T-joints, the dry preform was compacted into a T-shape acrylic fixture and analyzed via $\mu$CT, as shown in Figure 2. The web section of the preform was firstly fitted in the fixture

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**Figure 2.** $\mu$CT scan of the preform: (a) preform visualized via volume rendering from $\mu$CT showing the yarn directions; (b) schematic 3D orthogonal weave pattern (created by TexGen); (c) $\mu$CT section view of the preform with the numbers showing the numbering of weft layers.
before bifurcation of the flange section, thus no deformation in the web is likely to be expected from preform bifurcation and the part of the preform contacting the fixture corner would conform to its radius of curvature. The GE phoenix \textit{v}tomex m machine used a current of 240 $\mu$A and a voltage of 120 kV to achieve the imaging resolution of 30 $\mu$m per pixel in scanning. The compaction leads to an average fiber volume fraction of 45$\%$ which is calculated based on the preform areal weight with 4 mm thickness in the web and 2 mm in the flange. The preform has two geometric changes in its left-hand side: one is that the warp yarns at the junction (curved region) are in twice the filament count as those in the web; the other is a weft layer in the web (numbered as 4'(5') in Figure 2) is formed of twofold weft yarns as used in other layers and separated into two layers (layer 4' and 5') at the flange, for the purpose of yarn layer balance in the flange. Meanwhile, the flange part of the preform was woven with warp yarns of 24k filament count.

In order to validate the proposed geometric modeling approach, necessary geometric parameters measured from the preform are listed in Table 1. In addition, the intra-yarn fiber volume fractions at different locations inferred from the cross-sectional data were given, assuming a rectangular cross-section for the yarns. Measurements of each parameter were made using the software ImageJ at a number of $\mu$CT slices covering different yarns, from which the average and standard deviation (SD) were obtained. The folding process introduces additional yarn deformation at the junction, and three significant features were observed as described in the following sections.

### Warp yarn shift

The first distinctive deformation is the shift of the warp yarn stack at the junction, in the direction orthogonal to the yarn length direction. The warp yarns within the same stack are aligned vertically with each other before folding the preform, but their relative position is shifted after folding due to the rigid body transformation at the junction. However, the warp yarn shift could be suppressed by the surrounding binder yarns.

The shift begins in the warp yarn stack at the bend as marked in Figure 3, and is followed by the next 1–2 stacks in the flange. It is noted from the $\mu$CT scan that this phenomenon is most evident in the noodle area.
and the level of shift reduces in the following stacks of warp yarns in the flange. To quantify the shift of the warp yarn stack, shift in the bend can be represented by angle shift $S_{\theta}/C_1$, which is measured as the angle between the two lines connecting centers of yarns to the normal direction of each yarn. Shift of the following stack is depicted as displacement shift $S_d$ representing the horizontal distance between two yarn centers, as illustrated in Figure 3.

**Warp yarn cross-section bending**

Bent cross-sections of warp yarns were observed at the junction. As shown in Figure 3, each warp yarn bends about the centerline (or center-plane) of its cross-section uniformly along its length direction and its shape complies with the curvature of the adjacent weft yarn layers. For this 90° folded preform, the centerline of each weft layer in the bend follows the path of a quarter of the circumference of a circle. As a result, the centerlines of bent warp yarn cross-sections could be approximated with circular paths of different radii. The measured radii of curvature of centerlines of bent warp yarn cross-sections ($R_{\text{warp}}$) and weft yarn layers ($R_{\text{weft}}$) are listed in Table 2.

**Weft yarn cross-section flattening**

In Figure 4, section view A-A of Figure 2 shows the cross-sections of weft yarns at the middle of the bend arc, while section B-B denotes cross-sections of weft yarns in the flat region. It is obvious that the cross-sections of weft yarns at section A-A are significantly flattened, i.e. an increase in yarn width and decrease in yarn height, especially for those not sitting in the peak or trough of the binder yarn path since here yarn flattening is not constrained by binders. The cross-section data of flattened weft yarns without binder constraint at position A-A were measured from the µCT scan and results are listed in Table 1. The layer sequence of weft yarns in section A-A is numbered as Figure 2 shows. From further observation of µCT images, the most slender cross-section of weft yarn occurs at position A-A in Figure 2, from where it transitions symmetrically back to the dimension of the yarn cross-section in the flat part of the preform. It is also observed that the length of yarn width transition varies for yarns on different layers, but all the flattened yarns are of almost same height, as shown in Table 1.

Meanwhile, a number of gaps in yarn cross-sections at section A-A were observed, which caused serrated cross-section shapes in most of the weft yarns, as marked in Figure 4. A number of different slices near section A-A indicate this phenomenon and therefore it is unlikely to be an artefact of the µCT scan. The position of each gap appears random as observed and the cross-sectional gap is simply quantified by its maximum width that is named gap distance $G_d$. The value of $G_d$ is no more than 1.5 mm from several measurements. The flattening and cross-sectional gaps of weft yarns at the bend partly result from fiber migration when the preform is folded.15

**Table 2.** Measured and calculated radii of curvature of centerlines of bent warp yarn cross-sections and weft yarn layers, here subscript $n$ is the yarn layer number (mm; 5 measurements)

<table>
<thead>
<tr>
<th>Layer</th>
<th>$R_{\text{weft}}$</th>
<th>$R_{\text{warp}}$</th>
<th>$R_{\text{weft}}$</th>
<th>$R_{\text{warp}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Layer 1</td>
<td>2.19 (±0.07)</td>
<td>2.38 (±0.10)</td>
<td>2.18</td>
<td>2.4</td>
</tr>
<tr>
<td>Layer 2</td>
<td>2.42 (±0.12)</td>
<td>2.78 (±0.15)</td>
<td>2.62</td>
<td>2.84</td>
</tr>
<tr>
<td>Layer 3</td>
<td>3.11 (±0.08)</td>
<td>3.11 (±0.14)</td>
<td>3.06</td>
<td>3.28</td>
</tr>
<tr>
<td>Layer 4</td>
<td>3.63 (±0.12)</td>
<td>3.71 (±0.13)</td>
<td>3.5</td>
<td>3.72</td>
</tr>
<tr>
<td>Layer 5</td>
<td>4.25 (±0.10)</td>
<td>4.25 (±0.10)</td>
<td>3.94</td>
<td></td>
</tr>
</tbody>
</table>

**Geometric modeling**

**Warp yarn shifting**

The warp yarn shift is inevitable due to the rigid body transformation at the junction. Each warp yarn has a different radius of curvature of cross-sectional centerline to the fillet origin $O$. As illustrated in Figure 5, half of a flat piece of 3D woven preform (left) is folded to a 90° angle (right). $R_{\text{weft}}^{1}$, $R_{\text{warp}}^{1}$, and $R_{\text{weft}}^{2}$ are the radii of curvature of centerlines for weft yarns on layers 1, 2, and 3, respectively. As the radius of the inner surface of a T-piece is confined
by the mould corner radius \( R_m \), the following can be obtained:

\[
\begin{align*}
R_{\text{weft}}^1 &= R_m + \frac{H_{\text{weft}}}{2} \\
R_{\text{weft}}^2 &= R_{\text{weft}}^1 + D_{\text{weft}} \\
R_{\text{weft}}^3 &= R_{\text{weft}}^2 + D_{\text{weft}} \\
R_{\text{warp}}^n &= \frac{(R_{\text{weft}}^n + R_{\text{weft}}^{n+1})}{2}
\end{align*}
\]

where \( H_{\text{weft}} \) is the height of the weft yarn, \( D_{\text{weft}} \) is the spacing between through-thickness weft layers, and \( R_{\text{warp}}^n \) is the radius of curvature of centerline for the warp yarn cross-section on layer \( n \).

Assuming there is no yarn sliding, angle shift \( S_0 \) between warp yarn layer 1 and layer 2 (in radians) in Figure 5 can be expressed as

\[
S_0 = \frac{2(D_{\text{warp}} - d_o)}{R_{\text{weft}}^1 + R_{\text{weft}}^2} - \frac{2(D_{\text{warp}} - d_o)}{R_{\text{weft}}^2 + R_{\text{weft}}^3}
\]

\[
S_0 = (D_{\text{warp}} - d_o) \left( \frac{1}{R_{\text{warp}}^1} - \frac{1}{R_{\text{warp}}^2} \right)
\]

where \( D_{\text{warp}} \) is the warp yarn spacing within the same layer and \( d_o \) the offset distance from the centre of nearest unbent weft yarn describing the position where the weft yarns start to bend; therefore \( D_{\text{warp}} - d \) denotes the arc length in between bending onset position and bent yarn centre along the bend.

Also, the displacement shift \( S_d \) in adjacent layers for the following stacks of weft yarns can be described as

\[
S_d = \frac{\pi}{2} \left( \frac{R_{\text{weft}}^2 + R_{\text{weft}}^3}{2} - \frac{R_{\text{weft}}^1 + R_{\text{weft}}^2}{2} \right) = \frac{\pi}{2} D_{\text{weft}}
\]

\[\text{Warp yarn cross-section bending}\]

Yarn cross-section shapes have been approximated by a number of researchers by elliptical, power elliptical, and lenticular shapes which can be defined by parametric equations in two dimensions. Therefore, it is convenient to model the cross-section deformation by a parameterized transformation that can apply to every current parametric cross-section shape rather than defining a new description for a curved cross-section. The cross-section of one warp yarn is presented in Figure 6 prior to and after deformation, where \( x \) and \( y \) represent the original yarn width and yarn height.
directions, and the origin is at the yarn center. Accordingly the x and y coordinates of the original cross-section (the dashed line in Figure 6) can be defined by

\[
\begin{align*}
C(t)_x &= f(t) \quad 0 \leq t \leq 2\pi \\
C(t)_y &= g(t) \quad 0 \leq t \leq 2\pi
\end{align*}
\]  

(4)

It is assumed that the bent shape is a result of a simple conformation of the major axis from the axis \( x \) to the arc in radius \( R_{\text{warp}} \) (equation (1)). Thus, for each point on the upper half of the edge of a bent cross-section, the transformation can be achieved by adding \((\Delta x, \Delta y)\) to the coordinates of corresponding points on the original cross-section, where \( \Delta x, \Delta y \) can be defined by

\[
\begin{align*}
\Delta x &= (R_{\text{warp}} - H_{\text{warp}}/2) \sin \left( \frac{C(t)_x}{R_{\text{warp}} - H_{\text{warp}}/2} \right) - C(t)_x \\
\Delta y &= (R_{\text{warp}} - H_{\text{warp}}/2) \left( 1 - \cos \left( \frac{C(t)_x}{R_{\text{warp}} - H_{\text{warp}}/2} \right) \right)
\end{align*}
\]  

(5)

where \( H_{\text{warp}} \) is the height of the warp yarn. Similarly, for the point on the lower half of the edge of a bent cross-section, \( \Delta x, \Delta y \) are given by

\[
\begin{align*}
\Delta x &= (R_{\text{warp}} + H_{\text{warp}}/2) \sin \left( \frac{C(t)_x}{R_{\text{warp}} + H_{\text{warp}}/2} \right) - C(t)_x \\
\Delta y &= (R_{\text{warp}} + H_{\text{warp}}/2) \left( 1 - \cos \left( \frac{C(t)_x}{R_{\text{warp}} + H_{\text{warp}}/2} \right) \right)
\end{align*}
\]  

(6)

Usually the value of \( H_{\text{warp}}/2 \) is quite small compared with \( R_{\text{warp}} \), thus the above equations can be simplified into one transformation for all the points on the cross-section:

\[
\begin{align*}
\Delta x &= R_{\text{warp}} \sin \left( \frac{C(t)_x}{R_{\text{warp}}} \right) - C(t)_x \\
\Delta y &= R_{n} \left( 1 - \cos \left( \frac{C(t)_x}{R_{\text{warp}}} \right) \right)
\end{align*}
\]  

(7)

**Weft yarn cross-section flattening**

As shown in Figure 5 for half of a T-joint, each weft yarn is under a bending load. An idealized fiber distribution in the weft yarn would consist of parallel layers with concentric arcs whose arc lengths are dependent on each layer’s radial position. For illustration, Figure 7 represents a multi-filament yarn by inner, middle, and outer fibers/layers.

The model assumes that a fiber with high stiffness, such as carbon or glass fiber, has negligible axial deformation in a bent yarn. After folding a straight yarn (Figure 7(a)), it is clear that the assumed fiber arrangement in Figure 7(b) would lead to a fiber length difference which would not be realistic, in order to satisfy the boundary condition that a fiber bundle is undisturbed beyond the bending location. Meanwhile, Figure 7(c) shows another scenario where a fiber bundle with constant length complies with a 90° bend resulting in fiber ends at a sheared angle, which contradicts the real boundary condition. Therefore, to maintain constant fiber length at a localized 90° bend, fibers tend to migrate and rearrange resulting in a flattened yarn cross-section, as observed in Figure 4.

Taking the outermost layer of fibers in Figure 7(b) as a reference line for fiber movement, the rest of the fibers have to migrate to maintain their lengths. The positions of fibers except the reference ones in Figure 7(b) are considered as base positions for the movement. Hence the length differences between fiber on a specific fiber layer and fiber on the reference layer can be expressed as

\[
L_d = \frac{\pi}{2} \frac{d}{d}
\]  

(8)

where \( d \) is the distance from a specific fiber layer within the yarn to the reference layer. In practice, \( d \) can be a series of discrete values depending on the fiber arrangement within the yarn.

There are two possible mechanisms for a fiber to migrate that maintains fiber length at a bending corner. One mechanism is to move in the yarn height direction, referred to as in-plane migration (i.e. in the \( x-y \) plane shown in Figure 8). The other way is to move in the yarn width direction, referred to as out-of-plane
migration (z direction in Figure 9). In-plane fiber migration will result in a reduction of yarn height while out-of-plane migration will increase the yarn width.

**In-plane fiber migration.** A reference fiber and a fiber on a neighboring layer before and after migration are shown in Figure 8. The fiber loci in the bend can be approximated by circular paths as shown in Figure 3, thus in this model each single fiber path is described by one quadrant and two tangential straight lines, such as $R_0$, $l_{a0}$ and $l_{b0}$ in Figure 8 for example (where script $\theta$ or $d$ refers to reference or neighboring layer). The positions of the two tangent straight lines are assumed to be fixed in migration based on the $\mu$CT observation that weft yarn height is almost unchanged in the region away from the bend. According to the length relation in two axial directions shown in Figure 8, the following expressions can be obtained:

$$l_{a0} + R_0 = l_{a} + R_d + d \quad (9)$$

$$l_{b0} + R_0 = l_{b} + R_d + d \quad (10)$$

where $R_0$ is the radius of curvature of the reference fiber that can be expressed by the sum of radius of curvature of yarn centerline and half of the flattened yarn height; $R_d$ is radius of curvature of the fiber that is distance $d$ away from the reference fiber.

If a proportion $\alpha$ of the above defined arc length difference as shown in equation (8) is compensated by in-plane fiber migration, the arc length difference between migrated and base fibers is $\pi da/2$. Thus the relation between the length of the reference fiber and a migrated fiber on a layer distance $d$ away can be given as

$$l_{a} + \frac{\pi}{2} R_0 + l_{b} = l_{a} + \frac{\pi}{2} R_d + l_{b} + L_d (1 - \alpha) \quad (11)$$

**Out-of-plane fiber migration.** Since a proportion $\alpha$ of the defined fiber length difference is compensated by in-plane fiber rearrangement, the rest of the length difference, i.e. $(1 - \alpha)$ of the whole length difference, has to be compensated by out-of-plane fiber migration to meet

$$R_d = R_0 - d - \frac{\pi ad}{4 - \pi}$$

$$l_{a} = l_{a0} + \frac{\pi ad}{4 - \pi}$$

$$l_{b} = l_{b0} + \frac{\pi ad}{4 - \pi} \quad (12)$$

**Figure 7.** Idealized fiber arrangement in a 90° bent yarn.

**Figure 8.** Schematic view of in-plane fiber migration.

**Figure 9.** Schematic view of out-of-plane fiber migration.
the assumption that the folding process would not increase the arc length of each fiber within a 90° bent yarn. In theory fibers can migrate in both directions across the yarn width.

In Figure 9, $L_{ABCD}$ is the arc of a fiber which had in-plane but no out-of-plane migration so that its arc length is given by

$$L_{ABCD} = \overline{AB} + \overline{BC} + \overline{CD} = 2(\ell_0 - R_d) + \pi R_d$$  \hspace{1cm} (13)

where $\ell_0$ is the length of yarn width transition, denoted by length between A(or D) and G, from where the yarn width starts to change.

To calculate the arc length of a post-migration fiber $L_{AEFD}$, it is assumed that its radius of curvature is still the same as $L_{ABCD}$. $\overline{GH}$ denotes the distance that the fiber has migrated laterally, therefore in $\Delta ADH$ we have

$$\sin^2 \frac{\theta}{2} = \left(\frac{\overline{AD}/2}{\overline{AH}}\right)^2$$  \hspace{1cm} (14)

Also, in $\Delta AGH$ and $\Delta ADG$ we can have

$$\overline{AH}^2 = \overline{AG}^2 + \overline{GH}^2; \overline{AD}^2 = \overline{AG}^2 + \overline{GD}^2$$  \hspace{1cm} (15)

Hence, equation (14) is given by

$$\sin^2 \frac{\theta}{2} = \frac{1}{2} \frac{\ell_0^2}{l_0^2 + w^2}$$  \hspace{1cm} (16)

where for simplicity $l_0$ equals $\overline{AG}$ or $\overline{DG}$; $w$ equals $\overline{GH}$.

From trigonometry we can say that

$$\tan^2 \frac{\theta}{2} = \frac{l_0^2}{l_0^2 + 2w^2}$$  \hspace{1cm} (17)

Hence, in $\Delta EHO \overline{EH}$ is given by

$$\overline{EH} = \frac{\overline{EO}}{\tan \frac{\theta}{2}} = \frac{R_d}{\tan \frac{\theta}{2}}$$  \hspace{1cm} (18)

The arc length of migrated fiber $L_{AEFD}$ is

$$L_{AEFD} = \overline{AE} + \overline{EF} + \overline{FD} = 2\left(\sqrt{l_0^2 + w^2} - \frac{R_d\sqrt{l_0^2 + 2w^2}}{l_0}\right) + R_d\left(\pi - 2\tan\left(\frac{l_0}{l_0^2 + 2w^2}\right)\right)$$  \hspace{1cm} (19)

Thus, the relation between fiber arc length differences caused by out-of-plane fiber migration can be represented by

$$L_{AEFD} - L_{ABCD} = L_d(1 - \alpha)$$  \hspace{1cm} (20)

Substituting equations (13) and (19) into equation (20), the displacement for out-of-plane migration $w$ can be obtained for a fiber at a distance $d$ away from the reference fiber.

**Yarn cross-section.** The process of fiber migration resulting from bending is complex, as each single fiber may have a different value of $\alpha$ for its in-plane and out-of-plane migrations. Ideally, if all fibers on the same layer move with the same $\alpha$ and direction for out-of-plane migration, the cross-section of a weft yarn, assuming a square array for fiber packing, at Section A-A in Figure 2 would have a parallelogram shape as shown in Figure 10(b), where Figure 10(a) is the fiber arrangement before flattening with its left-hand layer as reference for in-plane migration (Figure 8). Since there are thousands of fibers in a single yarn, a large number of fibers presumably might migrate uniformly in the above way, i.e. without a change to their initial relative spatial sequence, as shown for the fibers represented by black dots in Figure 10(c) which remain in their ideal positions as in Figure 10(b). Meanwhile, a few fibers, not following the above ideal migration law, can penetrate into a position out of their own layers through more in-plane migration once out-of-plane migration of fibers on prior layers makes a space, as indicated by grey dots in Figure 10(c), which can cause a random serrated cross-section shape as observed in Figure 4. Additionally, a second possible cross-section shape is shown in Figure 10(d) if it is assumed fibers move with the same $\alpha$ but in both directions for out-of-plane migration. Hereafter the first large group of fibers following the pre-migration relative spatial sequence are named ideally migrated fibers, and the second small group are called randomly migrated fibers. Under this assumption, the values of $\alpha$ for the ideally migrated fibers are supposed to be identically denoted by $\bar{\alpha}$. Therefore the yarn height could be characterized by the distance between two corresponding fibers on the outermost layer (reference) and on the innermost layer, respectively, as shown in Figure 11.

If the flattened yarn height (at Section A-A in Figure 2) can be measured from a $\mu$CT scan, the value of $\bar{\alpha}$ for ideally migrated fibers on the innermost layer could be calculated by solving the following equation after substituting $R_{d_{\text{diss}}}$ with equation (12):

$$H_{wefl}^{\text{fl}} = R_0 - R_{d_{\text{diss}}} - \sqrt{2}dx$$  \hspace{1cm} (21)
where \( d_{\text{max}} \) is the maximum value of \( d \) which is the
original yarn height; \((-dx, -dy)\) in Figure 11 is the
position of the center of curvature of a fiber on the
innermost layer in the coordinate system of reference
fiber with its center of curvature as origin, which is
given by

\[
dx = dy = R_0 - R_{\text{dmax}} - d_{\text{max}}
\]

The yarn width is determined by the fibers that had out-
of-plane migration. Irrespective of randomly migrated
fibers and considering all fibers to be ideal ones migrating
in the yarn width direction with an identical \( \pi \), the relative
migration displacement \( w \) could be obtained by solving a
combination of equations (13), (19), and (20). In this case,
if the yarn extension, defined as the area beyond its ori-
ginal cross-section (area of fibers outside of the dotted
rectangle in Figure 10(b)), for the flattened yarn is full of
ideally migrated fibers, the total fiber area for fibers in the
yarn extension can be roughly estimated by

\[
\text{Area}_{\text{ext}} = \sum_{k=1}^{n_{f}} w_{k}d_{f}
\]

where \( d_{f} \) is the fiber diameter and \( n_{f} \) is the number of
total fiber layers along the yarn height direction before
flattening; \( w_{k} \) is the displacement of out-of-plane migra-
tion for a fiber on layer \( k \).

In reality, ideally migrated and randomly migrated
fibers coexist in a bent yarn so that a flattened cross-
section can be near to rectangular with the observed
gap (Figure 4) rather than a parallelogram (Figure 10
(b)) if all fibers are considered ideally migrated fibers.
Some of the fibers, namely randomly migrated fibers, in
the yarn extension do not have to move that much in
the yarn width direction because there is enough space
ahead of them to accommodate in-plane migration in
the yarn height direction, in the meantime the bending
force will prioritize in-plane migration. Therefore a
rectangular cross-section is likely to be formed, and a
rough prediction of weft yarn width extension \( W_{\text{ext}}^{\text{weft}} \) in
terms of constant fiber area in the yarn extension can be
made by

\[
W_{\text{ext}}^{\text{weft}} = \frac{\text{Area}_{\text{ext}}}{H_{\text{weft}}^{\text{fla}}}
\]

Figure 10. Schematic cross-section of fiber arrangement for a weft yarn showing fiber migration mechanisms: (a) pre-migration,
square packing fiber arrangement; (b) post-migration, idealized migration; (c) and (d) post-migration, possible realistic migration
accounting for serrated cross-section shape.

Figure 11. Yarn in-plane section view (\( \pi = 0.4; d_{\text{max}} = 0.35 \)).
where the value of $H_{\text{fla weft}}$ should be known in advance.

However, it is not possible that different fibers on the same layer can move in both directions simultaneously for out-of-plane migration as shown in Figure 10(d), because the measured $G_d$ is much smaller than double the value of maximum $w_k$ as calculated in the following section. Hence the predicted flattened yarn width can be expressed as

$$W_{\text{fla weft}} = W_{\text{weft}} + W_{\text{ext weft}}$$ (25)

where $W_{\text{weft}}$ is the original weft yarn width. This model is capable to predict the change of yarn cross-section if $H_{\text{fla weft}}$ can be determined.

**Validation**

Three proposed models for warp yarn shifting, cross-section bending, and weft yarn cross-section flattening were compared to experimental data from the standard orthogonal weave region of the preform as listed in Table 1.

**Warp yarn shifting**

The value of $R_m$ is 2 mm which equals the radius at the junction of the acrylic fixture contacting the preform. Based on equation (1), the radii of curvature of centerlines for warp cross-sections and weft yarns on different yarn layers can be obtained as listed in Table 2, which shows good agreement with the measured values.

The offset distance $d_0$ for the preform is measured as 2.0 mm and thus the relative angle shift of a warp yarn can be derived by equation (2). Furthermore, the displacement shift in adjacent layers can be determined by equation (3). Since the yarn shift is not significant for this preform, the angle shift measurements were only taken between yarns on layer 1 and yarns on other layers, whilst the displacement shift measurements were carried out between yarns on layer 2 and 3, as numbered in Figure 3. The warp yarns shift in a uniform manner but show a deviation along the yarn length direction. The measured shift values taken at different cross-sections along the length direction with standard deviations are shown in Table 3 in comparison with predicted values.

It is noticed that the measured angle shift between layer 1 and 3 is negative, which means the warp yarn on layer 3 has slid to the flange while folding. Also it is the same with displacement shift for yarns on layer 1 and 4 in the next warp stack.

Warp yarn stack shift can be observed but the trend does not always agree closely with the model since binder yarns influence the movement of warp yarns as well as yarn sliding induced by the weaving process.

**Warp yarn cross-section bending**

A bent cross-section shape can be obtained by the transformation defined in equation (7) on the basis of an original parameterized cross-section description. Figure 12(a) is based on an elliptical cross-section and its radius of curvature of centerline is 3 mm; Figure 12(b) is transformed from a super-ellipse (equation (26)) with $m = 0.4$ where radius of curvature of centerline of bent yarn cross-section is 2.4 mm.

It is observed that a real warp cross-section shape for this preform in the flat region can be better approximated with a super-ellipse than an ellipse, thus validation of warp yarn bending is based on the cross-section function for super-ellipse ($m = 0.4$) with yarn height and width measured from the unbent region in the web, where equation (4) can be expressed as

$$C(t) = \frac{W_{\text{warp}}}{2} \cos(t) \quad 0 \leq t \leq 2\pi$$

$$C(t) = \begin{cases} \frac{H_{\text{warp}}}{2} \sin^m(t) & 0 \leq t \leq \pi \\ -\frac{H_{\text{warp}}}{2} (-\sin(t))^m & -\pi \leq t \leq 0 \end{cases}$$ (26)

Then the bent cross-sections for different radii of curvature can be plotted and compared with real yarns, as shown in Figure 13. The modeled cross-section gives good agreement with realistic bent cross-section shapes.

**Weft yarn cross-section flattening**

Validation of cross-section flattening was carried out for yarns on three different layers in order to cover the full radius of curvature values which are the most essential parameters in yarn bending. The three layers

<table>
<thead>
<tr>
<th>Table 3. Warp yarn shift measured and predicted results (10 measurements)</th>
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<tbody>
<tr>
<td><strong>Angle shift $S_v$ (°) between</strong></td>
</tr>
<tr>
<td>---------------------------------------------------------------</td>
</tr>
<tr>
<td>layers 1 and 2</td>
</tr>
<tr>
<td>Measured value</td>
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<tr>
<td>Predicted value</td>
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selected are layers 1, 2, and 5 in section A–A of Figure 4. Weft yarn cross-section transition length $l_0$ is obtained from measurement for each layer, and radius of curvature of the reference fiber $R_0$ is derived based on calculated radius of curvature yarn centerline and flattened yarn height.

The total number of fiber layers along the yarn height direction before flattening $n_f$ (equation (23)) is estimated based on yarn dimensions before flattening and filament count as well as the assumption that the cross-section was a rectangular shape with square fiber packing. Based on the flattened yarn height, $\overline{d}$ can be obtained by solving a combination of equations (12), (21), and (22). Then the displacement of fiber out-of-plane migration along the yarn height if all the fibers are ideally migrated ones can be obtained by solving a combination of equations (13), (19), and (20), as plotted in Figure 14. Yarn extension width can be calculated from equation (23) and (24), and all the parameters for the prediction model are listed in Table 4.

Furthermore, the measured yarn width and predicted yarn width for the flattened weft yarns are compared in Figure 15 and reasonable agreement is observed.

Although the prediction accuracy is good, the model is semi-empirical based on the input parameters of flattened yarn height and cross-section transition length. However, it is interesting that all the flattened weft yarns with different radii of curvature have nearly the same height while the length over which cross-section transitions occur seems to increase linearly with radius of curvature in bending, both of which are merit further investigation.

**Figure 12.** Transformation of bent yarn cross-section shapes: dashed lines for original shapes and solid lines for bent shapes. (a) elliptical cross-section, (b) super-elliptical cross-section.

**Figure 13.** Model produced cross-section shapes in comparison with real yarns.

**Figure 14.** Relative displacement for fiber out-of-plane migration.
Conclusions

To characterize the relationship between the fiber architecture and composite mechanical performance, a numerical simulation requires a realistic geometric model as input. In this study, three geometric features, yarn shifting, cross-section bending, and flattening, caused by the bifurcation of a 3D flat woven preform for composite T-joints were identified by µCT analysis and approximated analytically. A simple model based on rigid body movement was used to approximate warp yarn shifting but the experimentally observed trend does not always agree closely with the model. The model would benefit from considering the influence of binder yarns on movement of warp yarns as well as yarn sliding induced by the weaving process but this would complicate the model significantly. Yarn cross-section bending was modeled by a transformation of the existing parameterized description of cross-section shape instead of proposing a new shape function, and good agreement was obtained when compared with µCT results. Lastly, weft yarn flattening was described in detail and a semi-empirical model was developed to predict the flattened yarn width and satisfactory agreement with experiments was observed. However, it is worthwhile to further study the preform compaction behavior and form a fully predictive model for yarn flattening.

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